

A note on Minimum Fisher information distributions with restricted support and fixed variance

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Abstract

Fisher information is of key importance in estimation theory. It also serves in inference problems as well as in the interpretation of many physical processes. The estimation error for the location parameter of a distribution is bounded by the inverse of the Fisher information associated with this distribution. In this note we look for minimum Fisher information distributions with a restricted support. More precisely, we study the problem of minimizing the Fisher information over the set of distributions with fixed variance defined on a bounded subset S of \mathbb{R} or on the positive real line. We show that the solutions of the underlying differential equation can be expressed in terms of Whittaker functions. Then, in the two cases considered, we derive the explicit expressions of the solutions and investigate their behavior. We also characterize the behavior of the minimum Fisher information as a function of the imposed variance.

Key words:

Fisher information, minimization problems, Whittaker functions, differential equations

1. Introduction

Importance of Fisher information as a measure of the information in a distribution is well known. It has many implications in estimation theory, as exemplified by the Cramér-Rao bound which is a fundamental limit on the variance of an estimator. It is used as a method of inference and understanding in statistical physics, as was promoted by Frieden, with many troublesome examples [7,9,8]. It is also used as a tool for characterizing complex signals or systems, [12,13,19] with applications, e.g. in geophysics [16,14,2], in biology [6], in reconstruction [3,4] or in signal processing [27,26,20]. In robust estimation, distributions minimizing Fisher information, in the case of scale and location parameters in a Kolmogorov neighborhood of a parent distribution, have been considered in [23,24]. Fisher information for orthogonal polynomials and special functions have been studied in [15,25]. Connections with the differential equations of Physics have been explored in [8].

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It is well known that the distribution with a fixed variance that minimizes the Fisher information on \mathbb{R} is the standard gaussian distribution. However, there are many situations where the variables at hands are known to belong to some subset of \mathbb{R} . For instance, the random variable may be known, on physical grounds, to have only non negative outcomes, e.g. the variable represents an energy. Variables may also be known to have a distribution with a support restricted to a given interval: this is the case of normalized variables or of the measurements obtained from a physical device with a (necessarily) finite output range. In such cases, it is of interest to consider, as a model of the data probability distribution, the distribution compatible with the constraints extracted from the data, with minimum (Fisher) information. This idea, advocated by Frieden and others, is reminiscent of the idea of maximum entropy distributions, when the measure of information considered is the Fisher information. By the Cramér-Rao bound, the Fisher information serves as a benchmark for estimators.

In the case of the estimation of a location parameter, the distributions with minimum Fisher information correspond to the most difficult estimation cases, and it is thus interesting to look for these distributions. This problem has been considered in the important paper [18], whose reading has tackled the present work. This paper presents general results and characterizations of the solution to the minimization of Fisher information on a compact support, subject to a variance constraint. Here, we give the explicit closed-form expressions of the solutions, characterize their behavior and propose alternate simpler proofs. We also extend the results to the case of distributions with support confined to the positive half line (and of course the case of any other semi-infinite support is a straightforward consequence).

Let f denote the probability density of a random variable X . The Fisher information (with respect to a translation parameter) is defined as

$$I[f] = \int_{\mathcal{S}} \left(\frac{d \ln f(x)}{dx} \right)^2 f(x) dx = \int_{\mathcal{S}} \left(\frac{d f(x)}{dx} \right)^2 \frac{1}{f(x)} dx, \quad (1)$$

where $\mathcal{S} = \text{Supp}[f]$ denotes the support of the density, $f(x)$ is supposed differentiable and both $f(x)$ and its derivative $f'(x)$ are square integrable on \mathcal{S} . We note \mathcal{D} the set of functions that verify these hypotheses. It is known [5] that the Fisher information is a strictly convex function of the distribution, that is: for $f(x)$, $g(x)$, and with $\lambda \in (0, 1)$, then

$$I[\lambda f + (1 - \lambda)g] < \lambda I[f] + (1 - \lambda)I[g]. \quad (2)$$

We are here interested in classifying distributions with a given variance, and consider the variational problem

$$I(\sigma^2) = \inf_f \{ I[f] : \text{Supp}[f] = \mathcal{S}, f \in \mathcal{D} \text{ and } \text{Var}[f] = \sigma^2 \} \quad (3)$$

which consists in finding a distribution with minimum Fisher information on the set of all distributions with support \mathcal{S} and a fixed variance σ^2 . The value of the minimum Fisher information obtained for a given variance σ^2 is denoted $I(\sigma^2)$ – the use of the square brackets and parenthesis distinguishes between the functionals of the probability distributions and the functionals of the variance. Although $I[f]$ is a convex functional, the set defined by the constraint $\text{Var}[f] = \sigma^2$ is not convex, so that uniqueness of the solution is not guaranteed. For instance, on \mathbb{R} , the normal distribution minimizes Fisher information in the set of distributions on \mathbb{R} with a given variance, but in fact irrespectively of the value of the mean: all normal distributions, whatever their mean, are equivalent solutions. We will obtain in the following that solutions on \mathbb{R}^+ or on an interval are in fact unique.

In section 2 we give the differential equation that is associated with the problem of minimization of Fisher information, and then we underline its relationships to known differential equations. So doing, we exhibit some explicit expressions of the solutions, in terms of Whittaker and parabolic cylinder functions. Then, we examine two particular cases. First, in section 4 we characterize the solutions with positive support, the behavior of the minimum Fisher information with respect to the variance, and show that the solution is unique and turns out to be a chi distribution. Second, in section 5 where the support is restricted to an interval, we give the expression of the unique solution to the problem and we study the general behavior of the minimum Fisher information with respect to the variance. Finally, we give the expression of the probability density with minimum Fisher information among all distributions with finite variance defined on an interval.

2. The differential equation associated with the minimization of Fisher information

In the following, it is convenient to introduce the transformation $f(x) = u(x)^2$ and to work with $u(x)$ instead of $f(x)$. With this notation, the Fisher information becomes

$$I[f] = 4 \int_{\mathcal{S}} u'(x)^2 dx, \quad (4)$$

where $u'(x)$ denotes the first order derivative of $u(x)$. The variational problem (3) can be restated as follows:

$$I(\sigma^2) = \inf_{u: u^2 \in \mathcal{D}} \left\{ 4 \int_{\mathcal{S}} u'(x)^2 dx : \int_{\mathcal{S}} u(x)^2 dx = 1 \text{ and } \text{Var}[u^2] = \sigma^2 \right\} \quad (5)$$

The problem above can also be completed by some additional conditions on the boundaries of the domain: for instance, when $\mathcal{S} = [0 : +\infty)$ we need $u(+\infty) = 0$ so as to ensure a proper integrable density. At the left endpoint, we shall ensure continuity and set $u(0) = 0$ in order to keep the Fisher information finite. Indeed, the Fisher information associated with a distribution is $+\infty$ when the distribution is not continuous on the domain \mathcal{S} . Similarly, when the support \mathcal{S} is restricted to $\mathcal{S} = [-1, 1]$, and the density is set to 0 outside of this domain, we shall take $u(1) = u(-1) = 0$ in order to ensure continuity at these boundary points.

The Lagrangian functional associated with the problem (5) is

$$L(u; \alpha, \beta) = \int_{\mathcal{S}} u'(x)^2 dx + \alpha \left(\int_{\mathcal{S}} u(x)^2 dx - 1 \right) + \beta \left(\int_{\mathcal{S}} (x - \mu)^2 u(x)^2 dx - \sigma^2 \right), \quad (6)$$

where α and β are the Lagrange parameters associated with the normalization and variance constraints respectively. The mean is denoted by μ . The minimum of the Lagrangian functional is obtained by standard calculus of variations [10] which asserts that a minimizer of (6) is necessary a solution of the Euler-Lagrange equation

$$-\frac{d}{dx} \frac{\partial L}{\partial u'} + \frac{\partial L}{\partial u} = 0 \quad (7)$$

which leads to

$$u''(x) - (\alpha + \beta(x - \mu)^2) u(x) = 0. \quad (8)$$

Of course, with the simple change of variable $z = x - \mu$, the differential equation reduces to

$$u''(z) - (\alpha + \beta z^2) u(z) = 0 \quad (9)$$

which is a parabolic differential equation. Interestingly, the minimum Fisher information can be written in terms of the constraints and Lagrange parameters.

Proposition 1 *The minimum Fisher information in (5) can be expressed as*

$$-\frac{1}{4}I(\sigma^2) = \alpha + \beta\sigma^2 \quad (10)$$

with $u(a) = u(b) = 0$ and where a and b denote the left and right endpoints of the support \mathcal{S} .

Proof. By integration by parts,

$$-\frac{1}{4}I(\sigma^2) = - \int_{\mathcal{S}} u'(x)^2 dx = - [u(x)u'(x)]_a^b + \int_a^b u(x)u''(x)dx. \quad (11)$$

Using the boundary conditions and the differential equation (8), we then obtain

$$-\frac{1}{4}I(\sigma^2) = \int_a^b (\alpha + \beta(x - \mu)^2) u^2(x)dx. \quad (12)$$

which reduces to (10) taking into account the values of the constraints. ■

3. Solutions to the differential equation

Let us consider the parabolic differential equation (9). Using the change of variable $x = \sqrt{2\sqrt{\beta}}z$, together with the substitution $d(x) = u\left((2\sqrt{\beta})^{-\frac{1}{2}}x\right)$, the differential equation becomes

$$d''(x) + \left(-\frac{\alpha}{2\sqrt{\beta}} - \frac{1}{4}x^2\right)d(x) = 0 \quad (13)$$

which is the Weber differential equation, whose standard form is

$$d''(x) + \left(\nu + \frac{1}{2} - \frac{1}{4}x^2\right)d(x) = 0. \quad (14)$$

Here we simply have $\nu = -a/2\sqrt{(b)} - 1/2$. The solutions of the Weber equation can be expressed as a linear combination of the parabolic cylinder function $D_\nu(x)$ and $D_{-\nu-1}(ix)$ [22]:

$$d(x) = c_1D_\nu(x) + c_2D_{-\nu-1}(ix) \quad (15)$$

But the Weber equation above can also be converted into the Whittaker equation using the substitution $d(x) = \frac{1}{\sqrt{x}}w\left(\frac{x^2}{2}\right)$ and $z = x^2/2$, which leads to

$$w''(z) + \left[\frac{3}{16z^2} + \frac{\nu + 1/2}{2z} - \frac{1}{4}\right]w(z) \quad (16)$$

that has the form of the Whittaker differential equation

$$w''(z) + \left[\frac{1/4 - \mu^2}{z^2} + \frac{\lambda}{z} - \frac{1}{4}\right]w(z) = 0 \quad (17)$$

with $\lambda = \nu/2 + 1/4$ and $\mu = 1/4$. But the Whittaker differential equation can also be obtained directly from the initial differential equation (9) with the substitution $u(z) = \frac{1}{\sqrt{z}}w\left(\frac{\xi z^2}{2}\right)$ and with $x = \xi z^2/2$. Then, one readily obtains

$$w''(x) + \left[\frac{3}{16x^2} - \frac{\alpha}{2\xi x} - \frac{\beta}{\xi^2}\right]w(x) = 0, \quad (18)$$

which reduces to the Whittaker differential equation (17) with $\xi = 2\sqrt{\beta}$, $\lambda = -\frac{\alpha}{4\sqrt{\beta}}$, and $\mu = \frac{1}{4}$. The general solutions of the Whittaker differential equation can be expressed in terms of the two linearly independent Whittaker functions $M_{\lambda,\mu}(z)$ and $M_{\lambda,-\mu}(z)$, or with the help of the Whittaker $W_{\lambda,\mu}(z)$ function defined by

$$W_{\lambda,\mu}(z) = \frac{\Gamma(-2\mu)}{\Gamma(\frac{1}{2} - \mu - \lambda)}M_{\lambda,\mu}(z) + \frac{\Gamma(2\mu)}{\Gamma(\frac{1}{2} + \mu + \lambda)}M_{\lambda,-\mu}(z) \quad (19)$$

for $2m \notin \mathbb{N}$. In these expressions, the Whittaker $M_{\lambda,\mu}(z)$ function can be expressed as a simple function of the confluent hypergeometric function, or Kummer function, $M\left(\frac{1}{2} + \lambda - \mu, 1 + 2\lambda, z\right)$ according to [1, eq. 13.1.32]

$$M_{\lambda,\mu}(z) = z^{\lambda+\frac{1}{2}}e^{-\frac{z}{2}}M\left(\frac{1}{2} + \lambda - \mu, 1 + 2\lambda, z\right). \quad (20)$$

Since the Kummer function has an exact series representation, we also have

$$M_{\lambda,\mu}(z) = z^{\mu+\frac{1}{2}}e^{-\frac{z}{2}} \sum \frac{\left(\frac{1}{2} + \lambda - \mu\right)_n}{n!(2\lambda + 1)_n} z^n, \quad (21)$$

where $(\)_n$ denotes the Pochhammer symbol. As far as the parabolic cylinder function is concerned, it can also be written as a Whittaker function [22, p. 347]:

$$D_\nu(x) = 2^{\frac{\nu}{2}+\frac{1}{4}}z^{-\frac{1}{2}}W_{\frac{\nu}{2}+\frac{1}{4},-\frac{1}{4}}\left(\frac{1}{2}z^2\right). \quad (22)$$

According to this discussion and the relationships between parabolic cylinder, Whittaker M and W functions, we find that the solutions of the differential equation associated with the problem of minimum Fisher information can be expressed as various equivalent linear combinations of the functions

$$D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}z) \text{ and } D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}iz),$$

or

$$\frac{1}{\sqrt{z}}M_{-\frac{\alpha}{4\sqrt{\beta}},\frac{1}{4}}(\sqrt{\beta}z^2), \frac{1}{\sqrt{z}}M_{-\frac{\alpha}{4\sqrt{\beta}},-\frac{1}{4}}(\sqrt{\beta}z^2), \text{ and } \frac{1}{\sqrt{z}}W_{-\frac{\alpha}{4\sqrt{\beta}},\frac{1}{4}}(\sqrt{\beta}z^2)$$

Some equivalent expressions of the solutions are given below

$$u(z) = c_1 D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}z) + c_2 D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}iz) \quad (23)$$

$$= c'_1 \frac{1}{\sqrt{z}}M_{-\frac{\alpha}{4\sqrt{\beta}},-\frac{1}{4}}(\sqrt{\beta}z^2) + c'_2 \frac{1}{\sqrt{z}}M_{-\frac{\alpha}{4\sqrt{\beta}},\frac{1}{4}}(\sqrt{\beta}z^2) \quad (24)$$

$$= c''_1 \frac{1}{\sqrt{z}}M_{-\frac{\alpha}{4\sqrt{\beta}},\frac{1}{4}}(\sqrt{\beta}z^2) + c''_2 \frac{1}{\sqrt{z}}W_{-\frac{\alpha}{4\sqrt{\beta}},\frac{1}{4}}(\sqrt{\beta}z^2). \quad (25)$$

In these formulas, the values of the constants in the linear combinations will be determined according to auxiliary constraints, e.g. the boundary values. We now turn to the characterization of solutions on \mathbb{R}^+ and then on an interval.

4. Solutions defined on the positive real line

We consider the expression of the solution in terms of the parabolic cylinder functions, as given in (23). A first point is to check that the solution is bounded on \mathbb{R}^+ . Since the Weber equation has only one irregular singularity at $z = +\infty$, it is sufficient to examine the behavior of the solution for $z \rightarrow +\infty$. The asymptotic expansion of the Weber function $D_\nu(z)$, with $\arg z \leq 3\pi/4$ is [22, p. 347], [1, eq. 19.8.1]:

$$D_\nu(z) \sim e^{-\frac{1}{4}z^2} z^\nu \left(1 - \frac{\nu(\nu-1)}{2z^2} + \frac{\nu(\nu-1)(\nu-2)(\nu-3)}{24z^4} - \dots \right). \quad (26)$$

For $z \rightarrow +\infty$, we see at once from the asymptotic expansion that $D_\nu(z) \rightarrow 0$. Replacing ν by $-\nu-1$ and z by iz , we can also observe that $D_{-\nu-1}(iz)$ tends to infinity when $z \rightarrow +\infty$. Therefore, since the solution must correspond to a proper integrable density, the second parabolic cylinder function in (23) must be discarded, with $c_2 = 0$, and the solution reduces to

$$u_{\alpha,\beta}(z) = \frac{D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}z)}{\left(\int_0^{+\infty} \left(D_{-\frac{\alpha}{2\sqrt{\beta}}-\frac{1}{2}}((2\sqrt{\beta})^{\frac{1}{2}}z) \right)^2 dz \right)^{\frac{1}{2}}} \quad (27)$$

where the denominator has been introduced in order to ensure the normalization of the probability density $f(z) = u(z)^2$. As a consequence of the simple form obtained above, it is possible to characterize the general behavior of the minimum Fisher information as a function of the variance. This is given by the following

Proposition 2 *The Fisher information of the minimum information probability density with positive support and fixed variance σ^2 , corresponding to the solution (27), verifies*

$$I_\nu(\sigma^2) = \frac{K_\nu}{\sigma^2}, \quad (28)$$

where $\nu = -\alpha/2\sqrt{\beta} - 1/2$ and K_ν is a constant.

Proof. Let $\nu = -\alpha/2\sqrt{\beta} - 1/2$ and $\xi = 2\sqrt{\beta}$, and consider ν fixed. Then, ξ acts as a scaling factor: more precisely, denoting $V_{\nu,\xi}$ the variance associated with the solution with parameters ν and ξ , we readily have $V_{\nu,\xi} = \frac{1}{\xi}V_{\nu,1}$ and $I_{\nu,\xi} = \xi I_{\nu,1}$, where $I_{\nu,\xi}$ is the Fisher information. Therefore, the product $I_{\nu,\xi}V_{\nu,\xi} = I_{\nu,1}V_{\nu,1}$ does not depend on ξ and

$$I_{\nu,\xi} = \frac{I_{\nu,1}V_{\nu,1}}{V_{\nu,\xi}}.$$

Moreover, from the equality $V_{\nu,\xi} = \frac{1}{\xi}V_{\nu,1}$ and the fact that $V_{\nu,1} > 0$, we deduce that the function $\xi \mapsto V_{\nu,\xi}$ maps \mathbb{R}^+ to \mathbb{R}^+ so that it is always possible to find a value ξ such that $V_{\nu,\xi} = \sigma^2$; thus, the Fisher information $I_\nu(\sigma^2)$ of the probability density associated with the solution $u(z)$ in (27) follows the equation (28), with $K_\nu = I_{\nu,1}V_{\nu,1}$. ■
Let ν^* be the value that minimizes $I_\nu(\sigma^2)$. For that value and any distribution f on \mathbb{R}^+ with same variance σ^2 , we always have

$$I[f] \geq I_{\nu^*}(\sigma^2) = \frac{K_{\nu^*}}{\sigma^2}, \quad (29)$$

which refines the Cramér-Rao inequality. We will check below that K_{ν^*} is of course greater than one. Actually, we know that the Fisher information is infinite in the case of a non differentiable density. It is thus important here to ensure continuity and differentiability at the origin. In order to ensure that the Fisher information of the probability distribution $f(x) = u^2(x)$ remains finite, we need to impose $u(0) = 0$. This implies a condition on ν so that $D_\nu(0) = 0$. It is easy to check that

$$D_\nu(0) = \frac{\sqrt{\pi}}{2^{\frac{\nu}{2}}\Gamma(\frac{1}{2} - \frac{1}{2}\nu)}. \quad (30)$$

Therefore, $D_\nu(0) = 0$ if and only if $\Gamma(\frac{1}{2} - \frac{1}{2}\nu) \rightarrow +\infty$, that is if and only if ν is an odd positive integer. In such a case, the Weber functions can also be expressed in terms of Hermite polynomials $H_n(x)$:

$$D_n(x) = 2^{-\frac{n}{2}}e^{-\frac{x^2}{4}}H_n(x). \quad (31)$$

Then, the determination of the optimum value ν^* of ν such that $I_\nu(\sigma^2)$ is minimum amounts to minimize K_ν with ν integer. This leads to the following result.

Proposition 3 *The distribution which minimizes the Fisher information on \mathbb{R}^+ , subject to a variance constraint, is obtained for $\nu^* = 1$, and is*

$$f(x) = \sqrt{\frac{2}{\pi}}x^2e^{-\frac{1}{2}x^2}, \quad (32)$$

the chi-distribution with three degrees of freedom. Its Fisher information, according to (45), is $I_1 = 3$, and its variance is $\sigma^2 = 3 - 8/\pi$. Then the minimum Fisher-variance product is $K_1 = 9 - 24/\pi \approx 1.3606$.

Proof. In the present case, it is possible to obtain closed-form formulas for the variance and information associated with (27), even for non integer values of ν .

The following relationships, which can be derived from integral representations of parabolic cylinder functions, see [21, 6.2] are useful:

$$\frac{d}{dx}D_\nu(x) = (-1/2)xD_\nu(x) + \nu D_{\nu-1}(x) \quad (33)$$

$$= (1/2)xD_\nu(x) - D_{\nu+1}(x) \quad (34)$$

Adding the two equalities and taking the square, we have

$$4\left(\frac{d}{dx}D_\nu(x)\right)^2 = \nu^2 D_{\nu-1}(x)^2 - 2\nu D_{\nu-1}(x)D_{\nu+1}(x) + D_{\nu+1}(x)^2. \quad (35)$$

Subtracting the equations (33) and (34), we obtain

$$xD_\nu(x) = D_{\nu+1}(x) - \nu D_{\nu-1}(x), \quad (36)$$

from which we deduce

$$xD_\nu(x)^2 = D_{\nu+1}(x)D_\nu(x) - \nu D_{\nu-1}(x)D_\nu(x), \quad (37)$$

$$x^2 D_\nu(x)^2 = D_{\nu+1}(x)^2 - 2\nu D_{\nu+1}(x)D_{\nu-1}(x) + \nu^2 D_{\nu-1}(x)^2, \quad (38)$$

The second ingredient of the calculation are the formulas [11, eqs. 7.711.2 and 7.711.3], from which we define the functional $S(\mu, \nu)$:

$$S(\mu, \nu) = \int_0^{+\infty} D_\mu(x)D_\nu(x)dx = \frac{\pi 2^{\mu+\nu+1}}{\mu - \nu} \left[\frac{1}{\Gamma(\frac{1}{2} - \frac{1}{2}\mu)\Gamma(-\frac{1}{2}\nu)} - \frac{1}{\Gamma(\frac{1}{2} - \frac{1}{2}\nu)\Gamma(-\frac{1}{2}\mu)} \right] \quad (39)$$

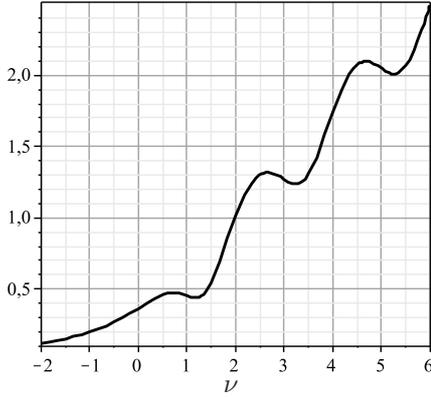


Fig. 1. Evolution of the variance $V_{\nu,1}$ in (44) with respect to ν .

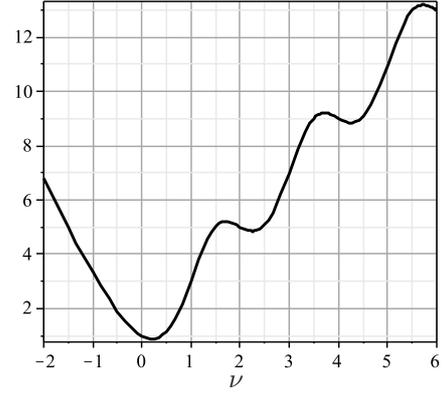


Fig. 2. Evolution of the information $I_{\nu,1}$ in (46) with respect to ν . Actually, the true Fisher information is only finite for positive odd values of ν .

for $\mu \neq \nu$, and with

$$S(\nu, \nu) = \int_0^{+\infty} D_\nu(x)^2 dx = \pi^{\frac{1}{2}} 2^{-\frac{3}{2}} \frac{\Psi\left(\frac{1}{2} - \frac{1}{2}\nu\right) - \Psi\left(-\frac{1}{2}\nu\right)}{\Gamma(-\nu)} \quad (40)$$

where $\Psi(x)$ is the digamma function. When μ or ν are integers, these formulas reduce to

$$S(2p, 2p + k) = -\frac{\pi 2^{2p + \frac{1}{2}k + \frac{1}{2}}}{k} \frac{1}{\Gamma\left(\frac{1}{2} - p\right)\Gamma\left(-p - \frac{1}{2}k\right)} \quad (41)$$

$$S(2p + 1, 2p + 1 + k) = \frac{\pi 2^{2p + \frac{1}{2}k + \frac{3}{2}}}{k} \frac{1}{\Gamma\left(-\frac{1}{2} - p\right)\Gamma\left(-p - \frac{1}{2}k\right)} \quad (42)$$

$$\text{and } S(m, m) = (2\pi)^{\frac{1}{2}} \frac{m!}{2} \quad (43)$$

So doing, using the expression of the solution (27), the definition of the Fisher information (4), equality (35) and the definitions (39-40), we obtain the expression

$$I_{\nu, \xi} = \xi \left(\nu^2 S(\nu - 1, \nu - 1) - 2\nu S(\nu - 1, \nu + 1) + S(\nu + 1, \nu + 1) \right) / S(\nu, \nu), \quad (44)$$

which reduces to the very simple expression

$$I_{m, \xi} = (2m + 1)\xi \quad (45)$$

in the integer case. Let us mention that a similar expression is reported in [25] for the case of Hermite functions, up to a factor 2 which is due to a different definition of Hermite functions.

In the same way, the integration of equalities (37-38) gives the first and second order moments, so the variance is

$$V_{\nu, \xi} = \frac{1}{\xi} \left(S(\nu + 1, \nu + 1) - 2\nu S(\nu + 1, \nu - 1) + \nu^2 S(\nu - 1, \nu - 1) \right) / S(\nu, \nu) - \left((S(\nu + 1, \nu) - \nu S(\nu - 1, \nu)) / S(\nu, \nu) \right)^2 \quad (46)$$

Figures 1 and 2 present the evolutions of the variance and information, as computed in (44) and (46). Figure 3 gives the information-variance product K_ν , as a function of ν . Actually, we know that the true Fisher information is only finite for positive odd values of ν (otherwise the density is discontinuous at the origin and the Fisher information is infinite).

Considering Figure 3, we read that the positive odd integer which minimizes the product $I_{\nu,1} \cdot V_{\nu,1}$ is $\nu^* = 1$. Accordingly, we obtain that the solution $f(x) = u(x)^2$ is (32). The values of its variance and Fisher information follow by direct computation. ■

Note that disregarding the differentiability requirement at the origin would lead to select the parameter $\nu = 0.1065$, corresponding to a variance $\sigma^2 = 0.38661$ and a ‘‘Fisher information’’ $I = 0.91886$. So doing one would obtain a

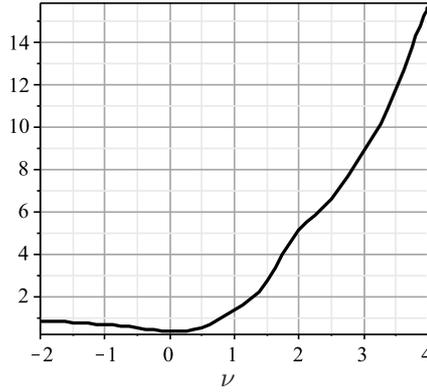


Fig. 3. The information-variance product K_ν as a function of ν .

product $I \cdot \sigma^2 = 0.35524$, which breaks the Cramer-Rao inequality. From the estimation theory point of view, it is clear that if the density has a bounded support and a discontinuity at the left endpoint of this support, then the variance of the estimate of the location parameter will be asymptotically zero, which corresponds to an infinite Fisher information in the Cramér-Rao bound.

5. Solutions with compact support

We know that the Fisher information associated with a distribution is invariant by translation of this distribution. Furthermore, a scaling of the distribution $f(x)$ according to $g(x) = \frac{1}{|a|}f(\frac{x}{a})$ with $a \neq 0$ yields a scaling of the Fisher information as $I[g] = \frac{1}{|a|^2}I[f]$. Hence, it is possible to restrict our study to any particular interval, and without loss of generality, we choose the interval $[-1, 1]$.

We consider the same problem as before, the minimization of the Fisher information subject to a variance constraint, and add the boundaries conditions $u(1) = u(-1) = 0$, which reads

$$I(\sigma^2) = \inf_{u: u^2 \in \mathcal{D}} \left\{ 4 \int_{-1}^1 u'(x)^2 dx : \int_{-1}^1 u(x)^2 dx = 1 \text{ with } Var(u^2) = \sigma^2 \text{ and } u(1) = u(-1) = 0 \right\} \quad (47)$$

The problem is invariant by the symmetry $x \mapsto -x$ since it is clear that if $u(x)$ is solution, then $v(x) = u(-x)$, which has the same variance and Fisher information, is also a solution. In fact, the distribution with minimum Fisher information for a given variance is unique and even. The general proof of this fact is rather involved and given in [17,18]. But since we have here the general expression of solutions to the underlying differential equation, it is not difficult to characterize this optimum solution.

Proposition 4 *The unique solution to the problem (47) is the non-negative function*

$$u(x) = \frac{\frac{1}{\sqrt{x}} M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}(\sqrt{\beta}x^2)}{\left(\int_{-1}^1 \frac{1}{\sqrt{x}} M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}(\sqrt{\beta}x^2) dx \right)^{\frac{1}{2}}}, \quad (48)$$

where $\sqrt{\beta}$ is the first zero of the function $M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}$.

Proof. We have already established that the general solution of the differential equation (9) can be expressed as a linear combination of two Whittaker M functions which are two linearly independent solutions:

$$u(x) = c_1 \frac{1}{\sqrt{x}} M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}(\sqrt{\beta}x^2) + c_2 \frac{1}{\sqrt{x}} M_{-\frac{\alpha}{4\sqrt{\beta}}, \frac{1}{4}}(\sqrt{\beta}x^2) \quad (49)$$

where c_1 and c_2 shall be chosen such that the boundaries and normalization conditions are satisfied. In this last equation, the first function is even while the second is odd, as it can be observed from the series development (21). Let us note them $S_{\text{even}}(x)$ and $S_{\text{odd}}(x)$ respectively. The boundaries conditions $u(1) = u(-1) = 0$ then imply that

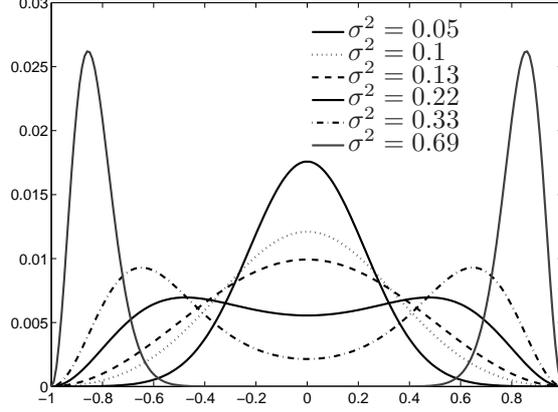


Fig. 4. Minimum Fisher information distributions with compact $(0, 1)$ support for several values of variances. For variances lower than $\sigma_*^2 = \frac{1}{3} - \frac{2}{\pi^2} \approx 0.1307$, the distributions are unimodal while they become bimodal when $\sigma^2 > \sigma_*^2$.

$$c_1 S_{\text{even}}(1) + c_2 S_{\text{odd}}(1) = 0 \quad (50)$$

$$c_1 S_{\text{even}}(1) - c_2 S_{\text{odd}}(1) = 0. \quad (51)$$

The only solution is $c_1 = c_2 = 0$, except if $S_{\text{odd}}(1)$ and $S_{\text{even}}(1)$ are simultaneously equal to zero. But it is easy to check that these two functions have no common zero. Therefore, the general solution can not include both terms. There shall be only one term, either $S_{\text{even}}(x)$ or $S_{\text{odd}}(x)$, in the general expression of the solution with its parameters adjusted so as to ensure the boundaries constraints.

We deduce that, since $u(x)$ is either odd or even, the associated density $f(x) = u^2(x)$ is even and has zero mean. In such a case, the set defined by the constraint $\text{Var}[f] = \sigma^2$ becomes a convex set and consequently the solution to the minimization of the Fisher information on this convex set is unique. From the fact both $u(x)$ and $v(x) = u(-x)$ are solution, we deduce that the unique solution is necessary the even one, that is (48).

Finally, if $u(x)$ is solution, so is $w(x) = |u(x)|$ because $w'(x)^2 = u'(x)^2$ and $w(x)^2 = u(x)^2$. By unicity of the solution we obtain that $u(x) \geq 0$ for $x \in (-1, 1)$. Therefore the condition $u(1) = 0$, the fact that all zeros of $M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}(x)$ are zero-crossings, and the non-negativity requirement, yield that the point $\sqrt{\beta}$ is the first zero of $M_{-\frac{\alpha}{4\sqrt{\beta}}, -\frac{1}{4}}$. ■

In Figure 4, we give some examples of solutions for several values $\sigma^2 \in [0, 1]$ of the variance. We see that for low variances the solution is unimodal while it is bimodal for higher variances. The Fisher information is large for low and high variances and much lower for intermediate variances. In fact, the minimum Fisher information is a convex function of the variance, cf [18, Theorem 6.1].

Proposition 5 *The minimum Fisher information $I(\sigma^2)$ is a strictly convex function of σ^2 .*

Proof. Let $u(x)^2$ and $v(x)^2$ be two distributions with minimum informations $I(\sigma_u^2)$ and $I(\sigma_v^2)$ respectively. Since the Fisher information is strictly convex, with $\epsilon \in (0, 1)$, we have

$$I[\epsilon u^2 + (1 - \epsilon)v^2] < \epsilon I[u^2] + (1 - \epsilon)I[v^2] = \epsilon I(\sigma_u^2) + (1 - \epsilon)I(\sigma_v^2) \quad (52)$$

Since the distributions u^2 and v^2 have zero mean, we have

$$\text{Var}[\epsilon u^2 + (1 - \epsilon)v^2] = \epsilon \sigma_u^2 + (1 - \epsilon)\sigma_v^2.$$

Then, there exists a distribution with the same variance and minimum Fisher information such that

$$I[\epsilon u^2 + (1 - \epsilon)v^2] \geq I(\epsilon \sigma_u^2 + (1 - \epsilon)\sigma_v^2). \quad (53)$$

Finally, combination of (52) and (53) yields

$$I(\epsilon \sigma_u^2 + (1 - \epsilon)\sigma_v^2) < \epsilon I(\sigma_u^2) + (1 - \epsilon)I(\sigma_v^2). \quad (54)$$

■

In the case of the \mathbb{R}^+ support, we obtained explicit expressions of both the variance and Fisher information, and as a result the exact behavior of $I(\sigma^2)$. The present case is more delicate. Indeed, although it is possible to obtain from (21) an exact series expansion of $u^2(x)$, where the coefficients depend of hypergeometric functions, the difficulty here is that we have no analytical expression for the argument of the first zero of the solution, which gives the value of β . Only approximations of this value are available [1, p. 510]. Therefore, we have to resort to a numerical determination of the parameters of the function. Similarly, we cannot give a close form formula for $I(\sigma^2)$. However, we can still characterize its general behavior.

Since $I(\sigma^2)$ is a convex function, it has a unique minimum, say σ_*^2 . The following result shows that this minimum is obtained for $\beta = 0$. Furthermore, this value discriminates two regimes for the solutions: in the case $\sigma < \sigma_*$, we have $\beta > 0$ while in the case $\sigma > \sigma_*$ we have $\beta < 0$ and the corresponding solution is the Whittaker M function with imaginary arguments. These facts have already been noticed in [18, Lemmas 5.2 and 5.3], but we provide here an alternate proof.

Proposition 6 *If σ_*^2 is the minimizer of $I(\sigma^2)$, then*

a) *for $\sigma < \sigma_*$, $\beta > 0$*

b) *for $\sigma > \sigma_*$, $\beta < 0$*

c) *and finally for $\sigma = \sigma_*$, $\beta = 0$*

Proof. Let $u(x)^2$ and $v(x)^2$ be two distributions with minimum informations $I(\sigma_u^2)$ and $I(\sigma_v^2)$ respectively, and let us define $g_\epsilon(x) = \epsilon u^2(x) + (1 - \epsilon)v^2(x)$. By Proposition 1, the Fisher information $I(\sigma^2)$ can be expressed as $I(\sigma^2) = -4(\alpha_v + \beta_v \sigma_v^2)$. We use the expansion proved in [17, Satz 7.2, p.90]:

$$I[g_\epsilon] = I[v^2] - 4\beta_v \epsilon (\sigma_u^2 - \sigma_v^2) + o(\epsilon^2). \quad (55)$$

Strict convexity gives

$$I[g_\epsilon] < \epsilon I[u^2] + (1 - \epsilon)I[v^2] = \epsilon I(\sigma_u^2) + (1 - \epsilon)I(\sigma_v^2) \quad (56)$$

Let us now take $\sigma_u = \sigma_*$. Since $I(\sigma_*^2)$ is the minimum Fisher information, we have the majorization

$$\epsilon I(\sigma_*^2) + (1 - \epsilon)I(\sigma_v^2) \leq I(\sigma_v^2), \quad (57)$$

and therefore $I[g_\epsilon] < I(\sigma_v^2)$. As a consequence, from (55) and the previous inequality, we obtain

$$\beta_v \epsilon (\sigma_*^2 - \sigma_v^2) > 0, \quad (58)$$

which gives cases a) and b) in the Proposition. The case c), $\beta = 0$, follows by continuity. ■

Hence, the solution which realizes the minimum $I(\sigma_*^2)$ of the Fisher information corresponds to $\beta = 0$. This means that this solution satisfies the differential equation $u''(x) = \alpha u(x)$, with $u(1) = u(-1) = 0$. This problem has for solutions $u(x) = \cos(k\pi x/2)$, with $\alpha = -k^2\pi^2$ and k integer. Since we know that the solution of the Fisher minimization is non-negative, there is only one possibility and

Proposition 7 *The probability density function $f(x) = u(x)^2$ defined on $[-1, 1]$ with minimum Fisher information is*

$$f(x) = \cos(\pi x/2)^2 \quad (59)$$

$$\text{with } \sigma_*^2 = \frac{1}{3} - \frac{2}{\pi^2} \text{ and } I(\sigma_*^2) = \pi^2. \quad (60)$$

Figure 5 reports the behavior of the Fisher information $I(\sigma^2)$. The Fisher information tends to infinity for $\sigma \rightarrow 0$ and $\sigma \rightarrow +\infty$ and its minimum is attained for $\sigma^2 = \sigma_*^2$. For $\sigma \rightarrow 0$, we obtain by [1, eqs. 13.1.32 and 13.5.5] that $u(x) \propto \exp(-\beta^{\frac{1}{2}}x^2/2)$, that is $u(x)$ converges to a Gaussian distribution with variance $\beta^{-\frac{1}{2}}$. Hence, for small values of the variance, the solution has the form of a Gaussian distribution concentrated on the origin, and, as a normal distribution, its Fisher information decreases as $1/\sigma^2$. When σ^2 increases towards 1, the two modes become more and more pronounced and probability accumulates on ± 1 , and the probability density function tends to two mass functions, as shown in Figure 4, and the Fisher information tends to infinity.

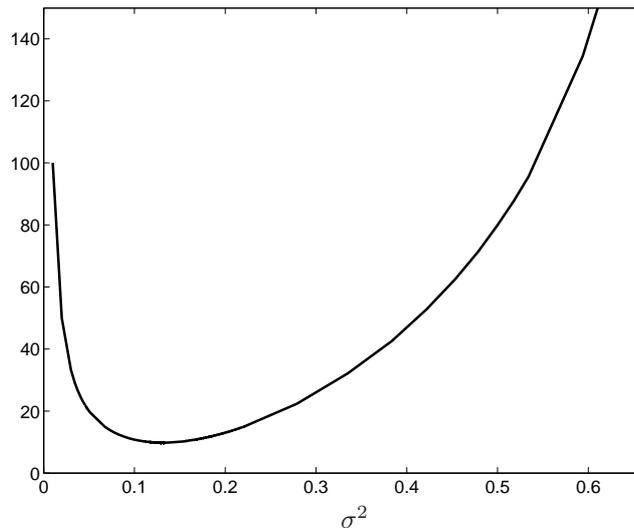


Fig. 5. Evolution of the minimum Fisher information $I(\sigma^2)$ in the case of a distribution with bounded support $(-1, 1)$

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